Evidence for Iodine Atoms as Intermediates in the Dye Sensitized Formation of I—I Bonds

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Photoinitiated reactions that form chemical bonds are of considerable importance for the conversion and storage of solar energy.1 While the O−O bond formation necessary for solar water splitting continues to be the “holy grail” in this area, there is also considerable interest in the formation of bonds between other elements such as halogens.1,2 For example, iodide oxidation to form I—I bonds is key to sensitizer regeneration in champion dye-sensitized solar cells. Mechanistic details of how iodide oxidation yields the I—I bonds present in I2− and I3− reaction products are under active investigation3–5 but remain speculative. Here we provide the first direct evidence that electron transfer sensitized to visible light with metal-to-ligand charge transfer (MLCT) excited states can directly yield iodine atoms that subsequently react to form an I—I chemical bond.

Figure 1 shows the absorption and photoluminescence (PL) spectra of [Ru(bpz)2(deeb)](PF6)2, where bpz is 2,2′-bipyrazine and deeb is 4,4′-(CO2Et)2-2,2′-bipyridine in neat acetonitrile. The measured PL quantum yield, φ, 0.14, and excited-state lifetime, τ = 1.75 ± 0.03 µs, were used to calculate the radiative and nonradiative rate constants, kφ = (8.36 ± 0.05) × 104 and knr = (5.01 ± 0.02) × 109 s−1. The absorption spectrum of the one electron reduced compound Ru(bpz)−(bpz)(deeb)+ generated by photochemical reduction is also shown in Figure 1. Ruthenium bipyrazine excited states are potent oxidants as previously reported6 and I3− formation of the MLCT excited state, Ru(bpz)2(deeb)2+ will be used to describe the excitation and nonradiative rate constants, kφ = (1.1 ± 0.1) × 105 M−1 s−1 from which kφ = 6.6 × 1010 M−1 s−1 was abstracted, Figure 2 inset. Steady-state PL measurements gave similar values.

Figure 2. Transient absorption difference spectra of Ru(bpz)2(deeb)+ measured at the indicated times after pulsed 532 nm excitation in acetonitrile solution. The inset shows excited-state decay measured as a function of the TBAI concentration with a corresponding Stern−Volmer plot from which Ksv = (1.1 ± 0.1) × 105 M−1 and kφ = 6.6 × 1010 M−1 s−1.

Pulsed laser excitation under conditions identical to those in Figure 2 except for the presence of 500 mM TBAI led to the appearance of new transient absorption features that were well described by equal concentrations of Ru(bpz)−(bpz)(deeb)+ and I2−. Figure 3. The 500 nm absorption band is due to the reduced ruthenium compound, and the UV and weak red absorption are mainly from I2−. Excited state quenching was quantitative, but the yield of charge separated products was low, φce < 0.1. Back electron transfer to yield ground-state products followed second-order equal concentration kinetics, kφ = (2.1 ± 0.3) × 1010 M−1 s−1.

At lower iodide concentrations (<1 mM), the appearance of the electron transfer products was quantified with nanosecond time resolution. Shown in the inset to Figure 3 are absorption transients measured at 400 nm as a function of the iodide concentration. This wavelength was chosen as it is an isosbestic point between the ground and excited state (Figure 2) and is very close to a ground-reduced state isosbestic point (λ = 404 nm, Figure 1). The formation and decay of I2− could therefore be quantified at this wavelength without interference from other species in the solution.7 A second-order rate constant for I2− formation was abstracted from iodide concentration dependent kinetic data like that shown in Figure 3A, kφ = (2.4 ± 0.2) × 1010 M−1 s−1.

The reduced ruthenium compound, Ru(bpz−)(bpz)(deeb)+, absorbs light most strongly at 500 nm. Kinetic measurements at this wavelength showed that the excited state decayed with the same state decay was found to be first-order with rate constants in good agreement with those abstracted from time-resolved PL data. The excited state was dynamically quenched by iodide with a Stern−Volmer constant Ksv = (1.1 ± 0.1) × 105 M−1 from which kφ = 6.6 × 1010 M−1 s−1 was abstracted, Figure 2 inset. Steady-state PL measurements gave similar values.

Figure 1. Room temperature absorption and photoluminescence spectra of Ru(bpz)2(deeb)+ in acetonitrile (black solid line). Also shown is the absorption spectrum of Ru(bpz−)(bpz)(deeb)+ (dashed red line).

Pulsed 532 nm excitation of [Ru(bpz)2(deeb)](PF6)2 in neat acetonitrile induces absorption changes that are consistent with formation of the MLCT excited state, Ru(bpz−)(bpz)(deeb)+. Positive absorptions at wavelengths <400 nm and >470 nm were assigned to the reduced bipyrazine ligand that is evident in Figures 1 and 2. Isosbestic points were observed at 400 and 540 nm. Excited

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rate constant as the appearance of the reduced compound. Second-order rate constants abstracted from iodide concentration dependent kinetic data yielded \((6.6 \pm 0.3) \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\) in excellent agreement with the value measured from Stern–Volmer analysis of the excited-state quenching, \((6.5 \pm 0.3) \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\).

**Scheme 1**

\[
\begin{align*}
\text{M}_\text{ox} + \Gamma & \rightarrow \text{M}_\text{red} + \Gamma \quad (1) \\
\text{M}_\text{ox} + 2\Gamma & \rightarrow \text{M}_\text{red} + \text{I}_2^- \quad (2)
\end{align*}
\]

The thermal oxidation of iodide by transition metal compounds has been proposed to occur by two distinct mechanisms, (1) and (2) in Scheme 1.\(^9\) It seems probable that excited-state reactions would follow similar pathways. Both reaction (1) and (2) are first-order in oxidized transition metal compound \(\text{M}_\text{ox}\) but are first- and second-order with respect to iodide. In principle, these pathways can be distinguished by identification of the reaction products \(\Gamma\) versus \(\text{I}_2^-\). In practice this is difficult as the iodine atom absorbs only weakly in the UV region and is known to rapidly react with iodide to generate \(\text{I}_2^-\). The second-order rate constant for \(\Gamma + \text{I}_2^- \rightarrow \text{I}_2^-\) in acetonitrile was reported to be \(2.3 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\).\(^8\) Therefore, observation of an \(\text{I}_2^-\) product is not sufficient proof for (2), as the first reaction can also ultimately yield \(\text{I}_2^-\).

Here we have shown that the rate constant for excited-state decay is, within experimental error, the same as that for the appearance of the reduced compound, which firmly establishes that excited-state electron transfer from iodide to the formally Ru(III) metal center in the excited-state produced \(\text{Ru(bpz)}^+(\text{bpz})\text{(deeb)}^+\) with a rate constant of \(k_{\text{eq}} = (6.6 \pm 0.3) \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\). The appearance of \(\text{I}_2^-\) was three times slower indicating that it is not an excited-state reaction product. Indeed, the measured \(k_v\) value is exactly that expected based on literature rate constants for the reaction of iodine atoms with iodide.\(^8\) Taken together this provides compelling evidence that the iodine atom is the initial product of excited-state electron transfer. The iodine atom subsequently reacts with iodide to form the \(\text{I}^-\text{I}^-\) bond of \(\text{I}_2^-\).

To our knowledge, this represents the first example of dye sensitized generation of iodine atoms. In previous work with \(\text{Ru(deeb)(bpy)}^2+\), an iodide ion was found to associate with each of the carbonyl ethyl ester oxygens in the solid state. The two iodides were 6.246 Å apart, and although it was not proven, the spectroscopic data in \(\text{CH}_2\text{Cl}_2\) were consistent with concerted excited-state electron transfer and \(\text{I}^-\text{I}^-\) bond formation.\(^10\) Related pathways have also been invoked at sensitized \(\text{TiO}_2\) interfaces where iodine pairing between the oxidized sensitizer and iodide is thought to facilitate the overall termolecular reaction (2).\(^11\) The results here show that iodine atoms are possible intermediates and that the rapid subsequent reaction with iodide in the 0.5 M LiI solutions typically used will significantly shorten their lifetimes.

**Scheme 2**

\[
\begin{align*}
E^* &= -0.82 \text{Ru}^{2+}(\text{bpy})^2+ \\
&\quad + \text{I}_2^-
\end{align*}
\]

\[
\begin{align*}
E^* &= -0.93 \text{Ru}^{2+}(\text{bpy})^2+ \\
&\quad + \text{I}^- + \Gamma
\end{align*}
\]

\[
\begin{align*}
E^* &= +1.36 \text{Ru}^{2+}(\text{bpy})^2+
\end{align*}
\]

The kinetic and mechanistic information reported in this communication along with iodide reduction potentials\(^12\) is summarized in the Jablonski-type diagram shown in Scheme 2. A bpz and deeb ligand was removed for clarity. Excited state electron transfer to yield the iodine atom is \(\sim 430 \text{ mV}\) downhill. Reaction of the iodine atom with iodide to make an \(\text{I}^-\text{I}^-\) bond lowers the free energy stored in the charge separated state by 110 mV.\(^9\)\(^12\) Charge recombination to yield ground-state products \(\text{Ru}^{2+} + \text{I}_2^- \rightarrow \text{Ru}^{3+} + 2\text{I}^-\) is highly thermodynamically favored (\(-\Delta G^\circ = 1.64 \text{ eV}\)) and occurs with a rate constant of \(2.1 \times 10^{10} \text{ M}^{-1} \text{ s}^{-1}\), almost 10 times larger than the \(\text{I}_2^-\) disproportionation rate constant. Unwanted charge recombination to \(\text{I}_2^-\) has previously been proposed to lower the efficiency of dye-sensitized solar cells,\(^2\) and these data show that it can be a very fast reaction.

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**References**


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