Investigation of Triphenylamine (TPA)-Based Metal Complexes and Their Application in Perovskite Solar Cells

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ABSTRACT: Triphenylamine-based metal complexes were designed and synthesized via coordination to Ni(II), Cu(II), and Zn(II) using their respective acetate salts as the starting materials. The resulting metal complexes exhibit more negative energy levels (vs vacuum) as compared to 2,2',7,7'-tetakis((N,N-di-p-methoxyphenylamine)-9,9′-spirobi fluorene (Spiro-OMeTAD), high hole extraction efficiency, but low hole mobilities and conductivities. Application of dopants typically used for Spiro-OMeTAD was not successful, indicating a more complicated mechanism of partial oxidation besides the redox potential. However, utilization as hole-transport material was successful, giving a highest efficiency of 11.1% under AM 1.5G solar illumination.

INTRODUCTION

In recent years, perovskite solar cells (PSCs) have attracted much attention due to their easy fabrication and high conversion efficiencies. Perovskites of this type are considered as the most promising materials to replace silicon solar cells in the market due to their high carrier mobilities, long charge diffusion length, high molar extinction coefficient, and strong absorption over almost the whole visible spectrum. However, at the same time, there are still some obstacles on the way toward commercialization, such as the toxic element lead, the stability of the solar cells, and the inclusion of the expensive hole-transport material (HTM) 2,2',7,7'-tetakis((N,N-di-p-methoxyphenylamine)-9,9′-spirobi fluorene (Spiro-OMeTAD). To avoid using lead in perovskite solar cells, Snaith and co-workers reported the CH3NH3SnI3 perovskite with an efficiency of 6.4% in 2014. Other structures, like CsSnI3 and (CH3NH3)(2)(BiI6),3 have also been studied and considered as potential lead-free light-absorbing materials. In the ambition to increase the stability of the solar cells, the traditional perovskite compound (CH3NH3PbI3) was modified to a mixed perovskite ((FAPbI3)0.85(MAPbBr3)0.15)20 or (FAI)0.81(PbI2)0.85-(MAPbBr3)0.15 representing one strategy. Other methods, such as passivation of the perovskite film,12,13 addition of hydrophobic interlayers,14 and encapsulation by polymers,15 have all been investigated. Also, much effort has been devoted to replacing the expensive Spiro-OMeTAD material. A previous report on the material X60 displays similar structure and comparable solar cell efficiency as cells containing Spiro-OMeTAD but a much simpler synthetic route, indicating a cost as low as U.S. $1/g. In comparison, the estimated corresponding Spiro-OMeTAD costs amount to more than U.S. $33/g.16

Metal complexes have shown applications in solar cells for many years. In the beginning, ruthenium(II) complexes were used as photosensitizers, like N3 and N719,17–19 in dyesensitized solar cells, being the basis for solar cells with very impressive performance. Later, cobalt20 and copper complexes21 were used as redox couples to replace the traditional I3−/I− system due to their tunable redox potentials and long-term stability. In recent years, metal complexes have also been used as hole-transport materials and dopants in perovskite solar cells, such as silver complexes,22 porphyrin,23 FK209,24 and copper complexes.25,26 However, the metal complexes reported as hole-transport materials have their limitations and cannot be widely up-scaled. In the silver system,22 the compound used is a silver complex cation with TFSI− as an anion offering a significant increase in conductivity, but the strategy was considered as doping. In the copper system,26 the structure of the solar cell was sandwich type, and by partial evaporation of the electrolyte, they could generate the so-called “Zombie” cells, which do not correspond to the traditional solid-state solar cells. Hence, we believe that it is meaningful to explore how a traditional metal complex works as a hole-transport material in all solid-state solar cell.
In this report, we use perovskite as our photosensitizer in all solid-state solar cell. At the same time, we designed and synthesized three different metal complexes with triphenylamine (TPA) derivatives as ligands coordinating to nickel, copper, and zinc, respectively. The aim is to investigate the performance of traditional metal complexes as HTMs and the influence from different metal ion centers on the resulting energy levels, conductivities, hole mobilities, and possible applications of these materials in perovskite solar cells. To the best of our knowledge, this is the first example that studies the effects from different coordinated metals as hole-transport materials in perovskite solar cells. Also, the first example when a nickel and a zinc complex have been investigated in combination with Spiro-OMeTAD.

### RESULTS AND DISCUSSION

The electron distribution of the highest occupied molecular orbitals (HOMOs) with HOMO-1s and the lowest unoccupied molecular orbitals (LUMOs) of the HTM molecules are shown in **Figure 1**. The HOMOs of these HTM molecules are mainly localized on the ligands. For Y1, the electron clouds are delocalized on both ligands, whereas in Y2 and Y3, the electron cloud is localized on predominantly one of the ligands, indicating that transport properties for holes may be different from that of Y1. All of the HOMO-1s have almost the same energy as HOMOs, but with electrons localized on the other ligand. Very little difference in energy between the HOMO and HOMO-1 confirms that the highest occupied MOs essentially are placed on the ligands. The LUMOs are all typically localized on the central metal atoms and adjacent pyrrole rings, and relative poor transport properties for electrons could be implied.

The redox potentials ($E_{\text{red}}$) of Y1, Y2, Y3, and Spiro-OMeTAD were determined using cyclic voltammetry (CV) in dichloromethane (DCM) solution (**Figure 2a**), collected with the corresponding data of energy levels versus vacuum from the calculations in **Table 1**. The HOMO energy level of Y1 is $-5.43$ eV (vacuum energy scale), which is a bit more negative than those of Y2 and Y3, both with an energy level of $-5.38$ eV, indicating that change in metal cores has only a slight effect on the redox potentials. This is in accordance with the theoretical calculations, which show that the electron density are mainly located in the peripheral parts of the ligands. Thus, it is the redox properties of the ligand that dominate the HOMO energy levels in these complexes and not the metal cores. It is worth mentioning that the redox peaks of Y1, Y2, and Y3 in the CV are irreversible, probably because of a demetallization process upon oxidation.\(^\text{27,28}\) The redox potential of Spiro-OMeTAD was determined to be $-5.12$ eV under the same condition, this is in good agreement with the previous report.\(^\text{29}\)

![Figure 1. Frontier orbitals of Y1, Y2, and Y3.](image)

![Figure 2. (a) Cyclic voltammograms of the metal complexes in dichloromethane. (b) UV−vis absorption spectra of the metal complexes in chloroform.](image)
reactions could be verified. First, we studied the hole-transfer properties at the perovskite-HTM interface using steady-state photoluminescence (PL) on a device with FTO/TiO2/perovskite/HTM. The results are shown in Figure 3. Devices with a naked perovskite film show a high light emission intensity due to the radiative recombination between photo-generated electrons and holes. When coating with HTMs, a significant decrease in the emission signals was observed for all of the materials. The PL quenching efficiencies for Y1, Y2, and Y3 are quite high and similar, most likely relating to their comparable HOMO energy levels and the corresponding driving force, indicating highly efficient hole transfer yields between the perovskite films and the HTMs, all above 90% and at the same level as Spiro-OMeTAD.

The reorganization energy is a key factor affecting the charge-transfer rate. It represents the structural relaxation upon the removal or addition of an electron to the system. Usually, molecular structures showing a smaller reorganization energy will also exhibit faster transfer rates for electrons or holes. As indicated in Table 2, Y2 shows the lowest calculated reorganization energy of 149 meV, which is comparable to 147 meV for Spiro-OMeTAD. However, compared with Y1 and Y3, the difference is minor; we should not expect a big difference in the carrier-transport ability among these materials.

Because it is always important to confirm the theoretical calculations from the experiment, we also studied their hole mobilities via the space-charge-limited currents method,31 see eq 1. Moreover, conductivities were determined using a two-contact electrical setup following a previously reported procedure,32,33 see eq 2. The results are shown in Figure 4.  

\[ J = \frac{9}{8} \mu_0 \epsilon_s \frac{V^2}{d} \]  

(1)

where \( J \) is the current density, \( \mu \) is the hole mobility, \( \epsilon_0 \) is the vacuum permittivity \( (8.85 \times 10^{-12} \text{ F m}^{-1}) \), \( \epsilon_s \) is the dielectric constant of the material (normally approximated to 3 for organic semiconductors), \( V \) is the applied bias voltage, and \( d \) is the film thickness.

\[ \sigma = \frac{L}{R} \mu d \]  

(2)

where \( \sigma \) is the conductivity, \( L \) is the channel length (about 2 mm), \( \mu \) is the channel width (about 10 mm), \( d \) is the film thickness (around 400 nm), \( R \) is the film resistance calculated from the gradients of the curves, and defined as the reciprocal of the slope, which is estimated from a least-squares fit of a linear model to the current—voltage data.

As shown in Figure 4a, pristine hole mobilities are quite low. All are around \( 2 \times 10^{-9} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1} \), as listed in Table 3. The differences are probably within the experimental errors. The
mobilities are lower than that of pristine Spiro-OMeTAD, which is about 1.69 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}. This is in good agreement with the small differences in calculated reorganization energies. Also, the new materials show lower conductivities as compared to Spiro-OMeTAD. For Y3, the conductivity is 4.05 \times 10^{-7} \text{ S cm}^{-1}. Y1 and Y2 are expected to display better carrier-transport properties. However, they are about the same as for Y3, approximately 4.06 \times 10^{-7} and 3.66 \times 10^{-7} \text{ S cm}^{-1}, respectively. Note that the values here are similar with the conductivity of the ligand (4) itself, which is 4.87 \times 10^{-7} \text{ S cm}^{-1} (Figure S10), showing that coordination does not have a positive effect on the increasing conductivities. Considering the small difference among Y1, Y2, and Y3, we believe that it is related to the structure of the metal complex itself. The designed molecular materials have very limited conjugated structures because the metal in the center breaks the whole molecule into two separate entities. Previous reports on pure organic molecules as HTMs have mentioned this effect, which emphasizes the importance of a large conjugated system. It realistic to assume that this also can be extended to metal complexes. On the other hand, it could also originate from the electron density distribution in the ground state. According to the Bässler model, the inherent dipolarity of a molecule has an adverse effect on the charge-carrier transport. In our systems, the HOMO electrons are localized on the ligands, which gives the ligands a negative charge and the central metal a positive charge. This charge separation causes a dipolarity in all of the metal complexes because the single electron pair in N in the triphenylamine unit could then easily transfer to N in the pyrrole ring, as shown in Figure 5.

Further, these materials were investigated as hole-transport materials in the perovskite solar cells. By decreasing the film thickness (lower to 50 nm), we could compensate for the disadvantage of low bulk charge transport properties. All three materials offered PSCs with voltages higher than 1 V, whereas a voltage of only 0.52 V could be obtained for solar cells without applied HTMs (Figure S11). However, the main problem of the new HTMs is still the low conductivity. As seen in Figure 6a, Y2 and Y3 do not significantly improve the photocurrent, showing 11.01 and 12.62 mA cm^{-2}, respectively. Thus, a poorer overall performance of 6.0 and 7.1% was obtained as compared to Spiro-OMeTAD (17.3%). Surprisingly, the best efficiency was obtained from the devices based on Y1, which offered a current density of 16.57 mA cm^{-2}, a voltage of 1.06 V, and a fill factor of 0.63 to yield an efficiency of 11.1%. The hysteresis phenomenon was also investigated and is shown in Figure 6a. All Y1, Y2, and Y3 display a high hysteresis with lower fill factor, indicating that metal complex hole-transport materials have minor effect on reducing the hysteresis. Corresponding data are summarized in Table 4. Distribution of 20 devices based on Y1, Y2, and Y3 as hole-transport materials in the

![Figure 4](image_url)

**Figure 4.** (a) J–V characteristics of thin films in hole-conducting-only devices without doping. (b) Current–voltage characteristics of different HTM films with same amount doping as for the solar cells.

<table>
<thead>
<tr>
<th>HTMs</th>
<th>hole mobility (cm² V⁻¹ s⁻¹)</th>
<th>conductivity (S cm⁻¹)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Y1</td>
<td>1.99 × 10⁻⁵</td>
<td>4.06 × 10⁻⁷</td>
</tr>
<tr>
<td>Y2</td>
<td>2.04 × 10⁻⁵</td>
<td>3.66 × 10⁻⁷</td>
</tr>
<tr>
<td>Y3</td>
<td>1.88 × 10⁻⁵</td>
<td>4.05 × 10⁻⁷</td>
</tr>
<tr>
<td>Spiro-OMeTAD</td>
<td>1.69 × 10⁻⁴</td>
<td>2.17 × 10⁻⁵</td>
</tr>
</tbody>
</table>

*a* Hole mobilities were estimated without doping. *b* Conductivities were determined with the same amount doping as for the solar cells.

![Figure 5](image_url)

**Figure 5.** Resonance structures of the metal complexes in the ground state.
perovskite solar cells are summarized in Figure S12, and a good reproducibility could be observed. The incident photon-to-current conversion efficiency (IPCE) spectra of PSCs with Y1, Y2, Y3, or Spiro-OMeTAD as the hole-transport material are shown in Figure 6b. The results clearly show that devices with Y1 have the highest IPCE compared with Y2 and Y3, although it is still lower than that of Spiro-OMeTAD. This is in good agreement with the \( J-V \) measurement, in which Y1 shows the highest current of 16.57 mA cm\(^{-2}\) among the three metal complexes. The integrated \( J_{sc} \) values from the IPCE spectra are 16.31, 10.65, 12.47, and 20.85 mA cm\(^{-2}\) for Y1, Y2, Y3, and Spiro-OMeTAD, respectively, matching the \( J-V \) test results quite well.

On the basis of the photovoltaic performance shown in Table 4, all of the new materials work as hole scavengers and perovskite solar cells are summarized in Figure S12, and a good reproducibility could be observed.

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<table>
<thead>
<tr>
<th>HTMs</th>
<th>scan direction</th>
<th>( J_{sc} )/mA cm(^{-2})</th>
<th>( V_{oc} )/V</th>
<th>FF</th>
<th>( \eta )/%</th>
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<tbody>
<tr>
<td>Y1</td>
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<tr>
<td>without HTM</td>
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<td>13.70</td>
<td>0.52</td>
<td>0.42</td>
<td>3.0</td>
</tr>
</tbody>
</table>

Figure 6. (a) \( J-V \) curves for devices based on Y1, Y2, Y3, or Spiro-OMeTAD as the hole-transport material. (b) Incident photon-to-current conversion efficiency (IPCE) spectra for devices based on Y1, Y2, Y3, or Spiro-OMeTAD as the hole-transport material.

Figure 7. SEM images of (a) the surface of the naked perovskite material, (b) the surface after coating with Spiro-OMeTAD, (c) cross section of a perovskite solar cell when Spiro-OMeTAD is applied as the hole-transport material, (d) the surface after coating with Y1, (e) the surface after coating with Y2, and (f) the surface after coating with Y3.
transport materials, but none is as good as Spiro-OMeTAD. The reason is mainly the low conductivities, as well as the low hole mobility, which hinder the charge-carrier transport. This limitation is also emphasized by an increase in the HTM layer thickness. Due to the low conductivity, an increase in layer thickness will result in a higher transport resistance, decreasing the current density and fill factor dramatically, as shown in Figure S13.

However, the film thickness can only be decreased to a certain thickness to compensate for this drawback. The scanning electron microscopy (SEM) images in Figure 7 show the film-forming properties of the different materials. Obviously, when using HTM concentration as low as 10 mg mL\(^{-1}\), the perovskite surface does not become well covered. Further decrease in thickness will cause an increase in the recombination losses because of uncovered perovskite coming in contact with the gold counter electrode material. Comparing the covered perovskite films by the metal complex materials when the same HTM concentration of the spin-coating solution was used, Y1 forms a much better film than Y2 and Y3. Figure 7e,f indicates that aggregation is a serious problem for the Y2 and Y3 systems. A large area with noncovered perovskite will allow direct contact with the gold counter electrode material, resulting in higher recombination losses and thus a smaller shunt resistance \((R_{sh})\) decreasing both the fill factor and the current as compared to Y1.

Furthermore, we investigated the impedance and capacitance using electrochemical impedance spectroscopy (EIS) of the solar cells based on Y1, Y2, Y3, or Spiro-OMeTAD as an HTM to study the interface properties. The results are shown in Figure 8. As expected, Y1 shows the lowest series resistance \((R_s)\), which is about 4 \(\Omega\), lower than those of the devices based on Y2 and Y3, which are 9 and 7 \(\Omega\), respectively. The lower series resistance is linked to a high PSC photocurrent. In addition, we note that although PSCs based on metal complexes show a lower recombination resistance \((R_{rec})\) than those with Spiro-OMeTAD, devices based on Y1 still show the highest \(R_{rec}\) among the three complexes, in good agreement with the photovoltaic performance in Figure 6a. We believe that this is because of the better film-forming ability of Y1 on the perovskite surface, preventing shunts between the perovskite and counter electrode materials. This effect is also indicated by the capacitances shown in Figure 8b. Devices based on Y1 display the lowest capacitance at low frequencies, whereas devices based on Y2 and Y3 show lower capacitances in the high-frequency region (>10\(^5\) Hz). This is because at low

![Figure 8](image_url)

**Figure 8.** (a) EIS impedance of solar cells based on Y1, Y2, Y3, or Spiro-OMeTAD as HTMs, (b) capacitance of solar cells based on Y1, Y2, or Y3 as HTMs.

![Figure 9](image_url)

**Figure 9.** Synthetic route to render the metal complexes.
frequency a better film of Y1 could efficiently extract and transport the holes. Thus, a lower capacitance could be achieved compared with Y2 and Y3. However, in the high-frequency region, the perovskite-HTM interface for devices based on Y1 experience a higher charge accumulation due to faster and better extraction ability, but a low carrier-transport ability. The excess of holes on the interface may result in a higher capacitance.

**CONCLUSIONS**

In summary, we have successfully designed and synthesized three metal complexes with TPA derivative ligands coordinated to nickel, copper, and zinc. These materials show very similar HOMO energy levels but an increase in band gaps going from nickel to copper to zinc. Also, these materials show high hole-extraction efficiencies but low hole mobilities and conductivities, resulting in satisfactory, but not excellent, PSC performance when applied as hole-transport materials. The best device performance was observed for the material based on nickel, with a conversion efficiency of 11.1%. This can most likely be ascribed the better film-forming ability on the perovskite surface. Application as dopants for Spiro-OMeTAD was not successful in spite of the more negative HOMO energy levels (vs vacuum). The current results will act as a guideline when exploring metal complexes as the hole-transport materials.

**EXPERIMENTAL SECTION**

Materials. Chemicals and solvents were purchased from Sigma-Aldrich and used without further purification except when mentioned. Lead iodide (PbI2, 99.99%) and lead bromide (PbBr2, 99.99%) were purchased from TCI and Alfa Aesar, respectively. Formamidinium iodide (FAI, CH(NH2)2I, >98%) and Ni(OAc)2·4H2O (74.6 mg, 0.3 mmol) were dissolved in 10 mL CH3OH and refluxed for 1 h. When the reaction finished, the mixture was concentrated by rotary evaporation and eluted with PE/EA = 2/1 to afford a light yellow solid (0.72 g, 80% yield).1H NMR (400 MHz, dimethyl sulfoxide (DMSO)-d6) δ 7.92 (d, 2H), 7.41 (d, 2H), 7.28 (d, 4H), 7.05 (d, 4H), 6.98 (d, 2H), 6.82 (d, 2H), 6.69 (d, 2H), 3.79 (s, 6H).13C NMR (100 MHz, DMSO-d6) δ 155.79, 147.63, 141.05, 133.92, 132.95, 128.89, 126.48, 120.73, 117.09, 114.70, 108.41, 107.04, 55.53, 43.32.

Chemicals and solvents were purchased from Sigma-Aldrich and distilled before use. 4-(Bis(4-methoxyphenyl)amino)benzaldehyde (2). 4-Methoxy-N-(4-methoxyphenyl)-N-phenylaniline (1). 4-Iodoanisole (6.5 g, 27.5 mmol), aniline (1 mL, 11 mM), and 1,10-phenanthroline (0.4 g, 2.2 mmol) were dissolved in 25 mL MeOH and refluxed for 1 h. When the reaction finished, the mixture was concentrated by rotary evaporation and eluted by recrystallization with DCM/MeOH, affording a dark red solid (88 mg, 62% yield).1H NMR (400 MHz, acetone-d6) δ 7.92 (d, 2H), 7.41 (d, 2H), 7.28 (d, 4H), 7.05 (d, 4H), 6.98 (d, 2H), 6.82 (d, 2H), 6.69 (d, 2H), 3.79 (s, 6H).

Ni-Di(4-((1H-pyrrol-2-yl)(2H-pyrrol-2-ylidene)methyl)-N,N-bis(4-methoxyphenyl)aniline) (Y1). 1H NMR (400 MHz, dimethyl sulfoxide (DMSO)-d6) δ 7.92 (d, 2H), 7.41 (d, 2H), 7.28 (d, 4H), 7.05 (d, 4H), 6.98 (d, 2H), 6.82 (d, 2H), 6.69 (d, 2H), 3.79 (s, 6H).13C NMR (100 MHz, DMSO-d6) δ 155.79, 147.63, 141.05, 133.92, 132.95, 128.89, 126.48, 120.73, 117.09, 114.70, 108.41, 107.04, 55.53, 43.32.

Cu-Di(4-((1H-pyrrol-2-yl)(2H-pyrrol-2-ylidene)methyl)-N,N-bis(4-methoxyphenyl)aniline) (Y2). 4-(Di(1H-pyrrole-2-yl)methyl)-N,N-bis(4-methoxyphenyl)aniline (3). 4-(Bis(4-methoxyphenyl)amino)benzaldehyde (2). 4-(Di(1H-pyrrole-2-yl)methyl)-N,N-bis(4-methoxyphenyl)aniline (3). 2 (1.6 g, 4.8 mmol) and freshly distilled pyrrole (40 mL, 0.44 mol) were treated with InCl3 (0.2 g, 1.0 mmol) at room temperature under nitrogen. After stirring overnight, the solution was filtered to remove InCl3. Excess pyrrole was further removed by rotary evaporation under reduced pressure. The residue was eluted by silica gel column with PE/EA = 2/1 to afford a light yellow solid (1.6 g, 86% yield).1H NMR (400 MHz, CDCl3) δ 7.97 (d, 2H), 7.12 (d, 4H), 6.90 (d, 4H), 6.88 (d, 2H), 3.82 (s, 6H). 4-(Di(1H-pyrrol-2-yl)methyl)-N,N-bis(4-methoxyphenyl)aniline (3). 2 (1.6 g, 4.8 mmol) and freshly distilled pyrrole (30 mL, 0.44 mol) were treated with InCl3 (0.2 g, 1.0 mmol) at room temperature under nitrogen. After stirring overnight, the solution was filtered to remove InCl3. Excess pyrrole was further removed by rotary evaporation under reduced pressure. The residue was eluted by silica gel column with PE/EA = 2/1 to afford a light yellow solid (1.6 g, 86% yield).
A (FA)\(_{0.85}(\text{MA})_{0.15}\)Pb(I\(_{0.85}\)Br\(_{0.15}\)\(_\text{perovskite film was made by spin-coating at 4500 rpm for 30 s inside an argon glovebox with a precursor solution containing 1.1 M PbI\(_2\), 1 M FAI, 0.2 M PbBr\(_2\), and 0.2 M MABr in a mixed solvent (DMF/ DMSO = 4/1). During spin-coating, 125 \(\mu\)L chlorobenzene was sprayed onto the film after 15 s. After finishing, the substrates were immediately placed on a hotplate at 100 °C for 30 min.

Hole-transport materials were spin-coated at 4000 rpm for 30 s using different precursor solutions. For Spiro-OMeTAD, the solution contained 70 mM Spiro-OMeTAD, 20 mM LiTFSI, 200 mM 4-tert-butylpyridine (TBP), and 2 mM FK209 in chlorobenzene. For materials Y1, Y2, and Y3, the concentration used were 10 mg mL\(^{-1}\), with 10 mM LiTFSI, 100 mM TBP, and 4% volume 1,1,2,2-tetrachloroethane (TeCA)\(^{37}\) as dopant (Figure S8). All of the solutions were prepared inside an argon-filled glovebox but spin-coated outside the glovebox.

The final step to fabricate a perovskite solar cell was the evaporation of a gold film as the counter electrode. A layer of 80 nm Au was deposited sequentially under high vacuum (<5 \times 10\(^{-5}\) Pa) by thermal evaporation (Edwards Auto 306) through a shadow mask to form an active area of around 20 mm\(^2\).

**Characterization.** NMR spectra were recorded on a Bruker AVANCE 400 MHz spectrometer. High-resolution mass spectra were collected with a Fourier transform-ion cyclotron resonance mass spectrometer instrument (Varian 7.0TFTICR-MS). Cyclic voltammetry (CV) was performed in dichloromethane with 0.1 M (TBA)PF\(_6\) as the supporting electrolyte, an Ag/AgNO\(_3\) electrode as the reference, a carbon-glass electrode as the working electrode, a Pt electrode as the counter electrode, and ferrocene/ferrocenium (Fc/Fc\(^+\)) as the internal reference using a CH Instruments electrochemical workstation (model 660 A). The UV–vis absorption spectra were recorded on a Lambda 750 spectrophotometer using a 1 cm cuvette. The concentration of the initial solution used was 10\(^{-3}\) M. The steady-state photoluminescence (PL) was investigated using a CARY Eclipse fluorescence spectrophotometer with an excitation wavelength of 530 nm. The scanning electron microscopy (SEM) images were recorded on a FEI-XL 30 series microscope. Film thickness was determined using a profilometer of model Veeco Dektak 150. The light source for the photocurrent–voltage (\(J–V\)) measurements was an AM 1.5G solar simulator (Newport 91160-1000). The incident light intensity was 100 mW cm\(^{-2}\) calibrated using a certified Si-photodiode (Fraunhofer ISE). The size of the working electrode was 25 × 3 mm\(^2\) on a substrate shared by four solar cells, and the counter electrode was 10 × 4 mm\(^2\). The solar cells investigated were masked to a working area of 0.126 cm\(^2\). The photocurrent–voltage (\(J–V\)) curves were obtained by the linear sweep voltammetry method using a Keithley 2400 source-measure unit with a scan rate 10 mV s\(^{-1}\). IPCE spectra were recorded on a computer-controlled setup with a light source ASB-XE-175, a monochromator (Spectral Products CM110), and a Keithley multimeter (model 2700). The setup was calibrated with a certified silicon solar cell (Fraunhofer ISE). Electrochemical impedance spectroscopy (EIS) and capacitance measurements were carried out using an Autolab PGStast12 potentiostat with an impedance module under a single sine function in dark with an applied bias potential of −1 V. The frequency range used was 0.1—100 000 Hz, together with an oscillating amplitude of 10 mV.
ASSOCIATED CONTENT

Supporting Information
The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsomega.7b01434.

1H NMR and 13C NMR of Y1 and Y3; mass spectra of Y1, Y2, and Y3; UV–vis spectra of Y3 with gradual addition of different amounts of TeCA as dopant; UV–vis absorption of Spiro-OMeTAD with gradual addition of different amounts of Y1, Y2, and Y3 as dopants; current–voltage characteristics of ligand (4) with doping; photovoltaic performance of perovskite solar cells without the hole-transport material; distribution of efficiencies of PSCs based on Y1, Y2, and Y3 as the hole-transport materials; photovoltaic performance with increased concentrations of Y1, Y2, and Y3 (Figures S1–S13) (PDF)

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Notes
The authors declare no competing financial interest.

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