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Non-Local Modelling of Strain Softening in Machining Simulations

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Abstract. Non-local damage model for strain softening in a machining simulation is presented in this paper. The coupled damage-plasticity model consists of a physically based dislocation density model and a damage model driven by plastic straining in combination with the stress state. The predicted chip serration is highly consistent with the measurement results.

1. Introduction

The introduction of strain or thermal softening to a material model gives rise to localized deformation. Local models need to be enriched to model the strain localization problems. An enrichment with a physical attraction is the non-local model, as discussed in previous studies [1]. Non-local models incorporate a characteristic length scale that can be related to the width of the localized zone. This enables a mesh-independent solution in the finite element analysis. However, its implementation in commercial codes requires simplification of the formulations due to the limitations in available user routines.

This work is concerned with the implementation of a non-local damage model in a commercial code, MSC.Marc[2] to be used for machining simulations. Thus, the ability to combine the implementation with remeshing is important. This approach is based on a previous study of the stability and accuracy of various simplifications required in the current context [3]. The plastic behaviour of the material is described by a physically based dislocation density plasticity model [4]. A damage model is coupled to this plasticity. Damage evolution takes into account the stress state and plastic strain of the material. In this paper, the simulation results show the effect of characteristic length scale in the damage model in terms of chip formation and its serrations.

2. Non-local damage modelling of strain softening

2.1. Damage – hypoelastic plasticity

Based on the local approach of fracture described in [5], the isotropic damage variable ω is considered to be the softening internal variable. Strain softening is then achieved by the use of effective stress, which is obtained as a transformation of the Cauchy stress σ as

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$$\bar{\sigma} = \frac{\sigma}{1 - h\omega} \tag{1}$$

where $0 \le h \le 1$ is a material parameter that differentiates between the damage behaviour in the tensile stress state and compressive stress state. In the case of tension, h = 1, Equation (1) reduces to its standard form as in [6].

The constitutive law of the undamaged material includes the damaged fraction by replacing the stress with the effective stress. This is possible because damage is considered to be the fraction of the material that does not transmit force/stress. The employed constitutive laws are summarized in Table 1, where C^e is the elastic material fourth order tensor. d^e and d^p are the elastic and plastic spatial velocity gradients, respectively. The right superscript ∇ denotes an objective stress rate. The effective stress given by Equation (1) is used in the yield criterion when determining the amount of plastic strain. Thus, $\overline{\sigma}_e$ is the von Mises stress of the effective stress tensor. σ_y is the yield limit for undamaged material. The caster \dot{A} is termed the plastic multiplier and corresponde to the rate of the

undamaged material. The scalar λ is termed the plastic multiplier and corresponds to the rate of the effective plastic strain in the current plasticity model.

Table 1. Constitutive model. (i) Additive split of the rate of deformation tensor $d = d^e + d^p$ (ii) Hypoelastic stress update $\sigma^{\nabla} = (1 - \omega)C^e : d^e$ (iii) Yield criterion $F = \overline{\sigma}_e - \sigma_y$ (iv) Plastic flow rule $\dot{\varepsilon}_e^p = \dot{\lambda} \frac{\partial F}{\partial \overline{\sigma}_e}$ (v) Rate of hardening variable $\dot{\varepsilon}_e^p = \dot{\lambda}$ (vi) Loading-unloading condition $\dot{\lambda} \ge 0; F \le 0; \dot{\lambda}F = 0$

2.2. Non-local damage model

The non-local formulation introduces a length scale by assuming that the state variables of the material constitutive equation not only depend on its local values at *x* but also depend on the values of one or several of the state variables in the domain around *x*. The size of this domain is the length scale parameter. The influence of this neighbourhood is accounted for by defining an integral of the non-local state variable v_{nl} , as introduced in a previous study [7] as

$$v_{nl}(\vec{x}) = \int_{\Omega(\vec{x})} \Phi(\vec{x}' - \vec{x}) v_l(\vec{x}') dx'$$
(2)

 v_l represents the local state variable in the continuous mechanics model. Φ is a weighting function, and $\Omega(\vec{x})$ is the material volume around the point \vec{x}_{\perp} in which Φ is not equal to zero.

It is assumed that the damage only evolves in the presence of plastic strain and tensile stress, as discussed in [8].

The rate of damage accumulation is given by

$$\dot{\omega} = \frac{\omega_c}{\varepsilon_r - \varepsilon_f} R_\theta R_\nu \dot{\varepsilon}_e^{p_+}$$
(3)

where $R_v = 2/3(1+v) + 3(1-2v)(\sigma_h/\sigma_e)^2$ includes the Cauchy stress triaxiality factor σ_h/σ_e .

 $\dot{\varepsilon}_e^{p+}$ is the accumulative plastic strain rate in tension, i.e., when $\sigma_h/\sigma_e \ge 0$, which is a simpler approach than splitting the stress tensor into positive and negative parts [9]. The term

$$R_{\theta} = \frac{1}{\left(\alpha_1 + \alpha_2 \overline{\theta}^2\right)} \tag{4}$$

is affected by the normalizedLode-angle $\overline{\theta}(-1 \le \overline{\theta} \le 1)$. This normalised angle in combination with the stress triaxiality characterise the stress state and is defined in [10] as

 $\overline{\theta} = 1 - \frac{2}{\pi} \arccos\left(\frac{27}{2} \frac{J_3}{\left(\sqrt{3J_2}\right)^3}\right)$ (5)

where J_2 and J_3 are the second and third invariant of σ . $\overline{\theta} = 1$ corresponds to axisymmetric tension, $\overline{\theta} = 0$ corresponds to plain strain or generalised shear loading and $\overline{\theta} = -1$ corresponds to axisymmetric compression [11]. The normalised lode angle has been found to characterize the influence of shear on damage [12]. In Equation (4), $\alpha_2 = (1 - \alpha_1)$. α_1 is an additional parameter to the original Lemaitre damage model [13]. Further variables in Equation (3) are \mathcal{E}_f , which is the effective plastic strain at damage initiation, and \mathcal{E}_r , which is the effective plastic strain at failure.

In non-local averaging, the internal variables are the best candidates when compared to other variables, such as the stress and total strain tensors [14]. The constitutive equation is formulated using non-local damage $\tilde{\omega}$. This is obtained by replacing the effective plastic strain in Equation (3) by its corresponding non-local variable $\tilde{\varepsilon}_{e}^{p}$ using Equation (2).

3. Experiments

The plastic deformation of AISI 316L stainless steel (see Table 2) at room and elevated temperatures was characterized using dynamic compression tests by means of a Split Hopkinson Pressure Bar (SHPB). The tests were performed with a maximum strain rate and temperature of 9000 s⁻¹ and 900 ${}^{0}C$, respectively. The stress-strain curves were used to calibrate the dislocation density model temperature dependent parameters. The parameters can be found in [15]. The plasticity model was used in the cutting simulations.

The geometry of the formed chips was obtained through quick-stop tests using equipment developed by SandvikCoromant. The chips show some serrations. The insert geometry is designed with an edge radius of 60 μm , a primary land of 0.15 mm and a chip rake angle of -6⁰. The cutting parameters consist of a cutting speed of 180 m/min and a feed of 0.15 mm/rev. The cutting depth was chosen to be 3 mm. This test case is used in finite element simulation in this paper, with the addition of strain softening.

Table 2. Chemical composition of AISI 316L (wt %).									
С	Si	Mn	Р	S	Cr	Мо	Ni	V	Ν
0.009	0.31	1.71	0.031	0.023	16.86	2.04	10.25	0.048	0.040

4. Finite element simulations

The cutting simulation is performed using the MSC.Marc [16]. The finite element simulation is modelled as a coupled mechanical and heat transfer analysis. The model is shown in **Figure 1**. The

workpiece length is 8 mm, and its height is 1.6 mm. The workpiece is meshed using a quadrilateral tool with triangular thermo-mechanical plane-strain elements. The tool is assumed to be thermoelastic. The initial model is composed of 900 elements in the workpiece and 2298 elements in the tool. The thermal and mechanical properties are given in [15], in which the plasticity model is described in detail.

The plasticity model was implemented in user routine WKSLP, whereas the non-local damage model was implemented in UDAMAG [16]. The separate implementation allows for additional robustness in cutting simulations, as different plasticity models can be coupled with different damage models. The parameters used for the Lemaitre model in Equation (3) are obtained from [17] for AISI 316L. The material parameters are presented in **Table 3**. The material constant h in Equation (1) is chosen as 0.001, so that nearly no damage accumulates during compression [8].

Chip formation is achieved via the continuous remeshing of the finite element mesh. There is no criterion for chip separation in the model. The remeshing is crucial for the stability of the simulation, as the analysis involves excessive deformations and strain localization in the cutting zone. The lowest bound for an element size in the workpiece is $2.5 \,\mu m$. This size is equivalent to the shortest tool element edge in the tool-chip contact interface. The maximum number of elements in the simulations is fixed at 20,000. To prevent severe element distortion, the mesh was updated every fourth increment. The element edge length determines the mesh refinement in the cutting simulations. The element edge length determines the mesh refinement in the cutting simulations. The meshes used in the simulations are described in

Table 4. The numerical solution applies an adaptive time stepping procedure. The time step is 10^{-6} seconds when chip formation has started and cutting forces become stable.



\mathcal{E}_{f}	Er	ω_c	α_1
0.035	0.502	0.395	1

Table 4. Finite element models showing the refinement variables used in the adaptive meshing.

Mesh	refinement	Refinement v	ariable,	Number of elements
notation		element edge length	(μm)	at steady state
Mesh1		100		12100
Mesh2		75		15200

5. Results

Figure 2 shows a comparison between simulations based on Mesh2 using local damage and non-local damage models. The result is very similar to that expected when the elements resolve at a length scale of $10 \,\mu m$. The element sizes are approximately $9 \,\mu m$. Refining the mesh will change the deformation of the local model as the localization is further enhanced by smaller elements. This notorious mesh dependency is remedied when using a non-local model [3,18].



Figure 2. Effective plastic strain rate using a fine mesh with (a) thermal and strain softening using a local damage model and (b) thermal and strain softening using a non-local damage model.

6. Conclusions

A non-local damage model of strain softening in machining simulations has been implemented into a commercial code, MSC.Marc [2].Chip serrations are predicted using the model. Explicit non-local updating [19] is used as it simplifies the implementation of the non-local damage model. The model also reduces the non-linearity of the non-local formulation and therefore makes it more robust at the cost of small time steps.

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